



Impact of relative permeability hysteresis on geological CO₂ storage

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[1] Relative permeabilities are the key descriptors in classical formulations of multiphase flow in porous media. Experimental evidence and an analysis of pore-scale physics demonstrate conclusively that relative permeabilities are not single functions of fluid saturations and that they display strong hysteresis effects. In this paper, we evaluate the relevance of relative permeability hysteresis when modeling geological CO₂ sequestration processes. Here we concentrate on CO₂ injection in saline aquifers. In this setting the CO₂ is the nonwetting phase, and capillary trapping of the CO₂ is an essential mechanism after the injection phase during the lateral and upward migration of the CO₂ plume. We demonstrate the importance of accounting for CO₂ trapping in the relative permeability model for predicting the distribution and mobility of CO₂ in the formation. We conclude that modeling of relative permeability hysteresis is required to assess accurately the amount of CO₂ that is immobilized by capillary trapping and therefore is not available to leak. We also demonstrate how the mechanism of capillary trapping can be exploited (e.g., by controlling the injection rate or alternating water and CO₂ injection) to improve the overall effectiveness of the injection project.

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1. Introduction

[2] There is little doubt that human socioeconomic activities are interacting with the biogeochemistry of the planet at a global scale. One example is the impact of carbon dioxide atmospheric emissions on the global carbon cycle [Falkowski et al., 2000]. The result of anthropogenic releases is an accumulation of carbon dioxide in the atmosphere, accompanied by a reduction in the pH of the upper ocean [Intergovernmental Panel on Climate Change, 1996]. It is also well documented that carbon dioxide is a greenhouse gas, and one of the main contributors to global warming [Falkowski et al., 2000; Cox et al., 2000]. Current predictions suggest that, unless an aggressive reduction of net CO₂ emissions is implemented, carbon dioxide concentrations in the atmosphere will continue to rise during this century [Wigley et al., 1996; Hoffert et al., 1998]. Since anthropogenic CO₂ emissions are primarily due to energy consumption, and 85% of the primary power is supplied by fossil fuels now, a drastic reduction in CO₂ emissions represents a major challenge [Orr, 2004].

[3] CO₂ sequestration refers to the capture and long-term storage of anthropogenic CO₂ in order to limit its emission to the atmosphere [Lackner, 2003]. Injection into geological

formations is one option to store CO₂ [Hitchon et al., 1999; Bachu, 2000; Orr, 2004]. Different target formations have been identified for this purpose, including depleted oil and gas reservoirs [Holloway, 2001; Kovscek and Wang, 2005; Kovscek and Cakici, 2005], unminable coal beds [Bromhal et al., 2005], and deep saline aquifers [Bruant et al., 2002; Pruess and Garcia, 2002; Pruess et al., 2003; Bachu, 2003].

[4] One of the major concerns in any sequestration project is the potential leakage of the CO₂ into the atmosphere. Possible causes of leaks are loss of integrity of the cap rock due to overpressurization of the geological formation [Rutqvist and Tsang, 2002; Jimenez and Chalaturnyk, 2002], and abandoned wells that may be present [Nordbotten et al., 2004]. When planning geologic sequestration projects in saline aquifers or depleted hydrocarbon reservoirs, it is therefore essential to predict the migration and distribution of the CO₂ in the subsurface structure so that injection can be maximized while keeping the risk of leakage at a minimum.

[5] In this investigation we consider injection and storage of CO₂ in saline aquifers. Many authors have presented simulations of CO₂ injection and migration [see, e.g., Johnson et al., 2000; Ennis-King and Paterson, 2002; Wellman et al., 2003; Pruess et al., 2003; Xu et al., 2003; Doughty and Pruess, 2004; Flett et al., 2004; Kumar et al., 2005; Mo and Akervoll, 2005; Obi and Blunt, 2006] using a variety of approaches. Because of the density difference between the CO₂ and the brine, the low viscosity CO₂ tends to migrate to the top of the geologic structure. This upward migration is sometimes delayed or suppressed by low permeability layers that impede the vertical flow of gas (in this paper we often refer to the CO₂ phase as gas even

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though, in practice, it is typically present as a supercritical fluid). There are several mechanisms by which the CO₂ can be stored, which include the following. (1) In hydrodynamic trapping, the buoyant CO₂ remains as a mobile fluid but is prevented from flowing back to the surface by an impermeable cap rock [Bachu *et al.*, 1994]. (2) In solution trapping, dissolution of the CO₂ in the brine [Pruess and García, 2002], possibly enhanced by gravity instabilities due to the larger density of the brine–CO₂ liquid mixture [Ennis-King and Paterson, 2005; Ennis-King *et al.*, 2005; Riaz *et al.*, 2006]. (3) In mineral trapping, geochemical binding to the rock due to mineral precipitation [Gunter *et al.*, 1997; Pruess *et al.*, 2003]. Finally, (4) in capillary trapping, disconnection of the CO₂ phase into an immobile (trapped) fraction [Kumar *et al.*, 2005; Flett *et al.*, 2004; Spiteri *et al.*, 2005].

[6] In this paper we demonstrate that the picture that emerges from the first three sequestration mechanisms, permeability, dissolution and mineral trapping, is incomplete. What is missing is consideration of the irreversibility of multiphase flow dynamics and, in particular, relative permeability hysteresis and physical trapping of the CO₂ phase within the porous medium. During the injection period, the less wetting CO₂ displaces the more wetting brine in a drainage-like process. However, after injection, the buoyant CO₂ migrates laterally and upward, and water displaces CO₂ at the trailing edge of the plume in an imbibition-like process. This leads to disconnection of the once-continuous plume into blobs and ganglia, which are effectively immobile [Hunt *et al.*, 1988]. It is this sequestration mechanism, capillary or residual trapping of CO₂ in its own immobile phase, that we emphasize in this investigation. We show that the capillary trapping mechanism has a huge impact on the migration and distribution of CO₂ which, in turn, affects the effectiveness of the other sequestration mechanisms.

[7] The importance of the “residual” CO₂ saturation has been pointed out by Doughty and Pruess [2004] and Hovorka *et al.* [2004], with reference to laboratory and field data from the Frio brine pilot experiment. However, no distinction was made between critical saturation (during drainage) and residual saturation (during imbibition) in their simulations. Kumar *et al.* [2005] performed a study of CO₂ storage in saline aquifers that accounted for dissolution and chemical reaction. They considered relative permeability hysteresis using a Land-type model. They concluded that the effect of residual gas on CO₂ storage can be very large and more significant than sequestration as brine or mineral. Their results indicated that significant capillary trapping of CO₂ occurred during the injection phase. However, injection is a drainage process, and gas is not trapped during drainage: Residual gas is formed after injection stops when water displaces CO₂. Postinjection trapping was analyzed in the recent study by Flett *et al.* [2004], who used a Land hysteresis model and emphasized the sensitivity of the predictions to the choice of the Land trapping parameter. In our simulations we will confirm that trapping is a significant storage process, but only after the initial injection phase.

[8] There seems to be some consensus that the timescales associated with each of the sequestration mechanisms are very different [Pruess *et al.*, 2003]. Residual and hydrody-

amic trapping occur on small timescales, dictated by the two-phase flow dynamics. Despite the fact that dissolution into the aqueous phase can be considered as an equilibrium process locally, it can be slow at the aquifer scale because it relies on the diffusion of CO₂ to regions not in contact with the CO₂ plume. Unstable gravity-driven convection can increase the rate of dissolution (because brine saturated with CO₂ is slightly more dense than brine alone), but the estimates of Ennis-King and Paterson [2005] and Riaz *et al.* [2006] indicate that the increase is significant only in relatively high permeability formations. Mineral trapping is believed to be a very slow process, because of the slow kinetics of precipitation reactions [Knauss *et al.*, 2005]. Therefore it is not unlikely that the timescale for mineral trapping is at least an order of magnitude larger than that of solution trapping, which may easily be an order of magnitude larger than that of hydrodynamic and residual trapping. As a result, we investigate the effects of capillary and hydrodynamic trapping alone, and we establish that relative permeability hysteresis becomes a major factor in the assessment of CO₂ sequestration projects. Moreover, we elucidate how the irreversible character of two-phase flow can be exploited to maximize residual trapping and subsequently increase CO₂ storage.

[9] An outline of the paper is as follows. In section 2 we give an account of the phenomenon of hysteresis from a pore-scale viewpoint, and explain what is the physical basis for residual trapping in aquifer disposal of CO₂. In section 3 we describe the setup of the numerical simulations employed to assess the importance of hysteresis. We used a realistic three-dimensional heterogeneous formation, as well as measured relative permeability data that display hysteresis. The results of the simulations are discussed in section 4. Finally, in section 5, we gather the main conclusions of this investigation.

2. Trapping and Relative Permeability Hysteresis

[10] Hysteresis refers to irreversibility, or path dependence. In multiphase flow, it manifests itself through the dependence of the relative permeabilities and capillary pressures on the saturation path and the saturation history. From the point of view of pore-scale processes, hysteresis has at least two sources [de Gennes *et al.*, 2004]. (1) The first source is contact angle hysteresis: the advancing contact angle (of wetting phase displacing a nonwetting phase) is larger than the receding contact angle (of wetting phase retreating by nonwetting phase invasion) due to chemical heterogeneities or surface roughness. (2) The second source is trapping of the nonwetting phase: during an imbibition process, a fraction of the nonwetting phase gets disconnected in the form of blobs or ganglia, becoming effectively immobile (trapped). Hysteresis effects are larger in processes with strong flow reversals. This is the case of cyclic water and gas injection in a porous medium, in which the gas phase is trapped during water injection after a gas flood.

2.1. Pore-Scale View of the Hysteresis Phenomenon

[11] Trapping and hysteresis have been successfully explained in terms of the displacement mechanisms that take place at the pore scale [Lenormand *et al.*, 1983], and successfully modeled using pore-network simulation tools

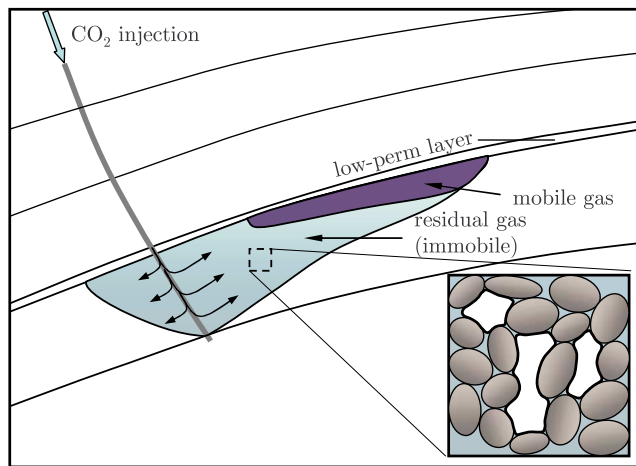


Figure 1. Schematic of the trail of residual CO₂ that is left behind because of snap-off as the plume migrates upward during the postinjection period.

[e.g., *Jerauld and Salter, 1990; Valvatne and Blunt, 2004*]. Consider a medium that is initially filled with water. The solid grains are made of minerals that are naturally wetting to water and, therefore the medium is water wet. During CO₂ injection into the aquifer, the nonwetting CO₂ phase invades the pore space. This is a drainage process in which the only mechanism for displacement of water by CO₂ is piston-type displacement: the CO₂ invades the porous medium in the form of a continuous, connected cluster. Water, however, remains present not only in small pores that have not been filled with CO₂ but also in the corners and crevices of the pores that have been invaded. Consider now the displacement of the CO₂ by water. During this process, there are several physical mechanisms by which the water can displace the CO₂ [*Lenormand et al., 1983*]. In addition to piston-type displacement, cooperative pore-body filling and snap-off may occur. For water wet rocks, snap-off is the dominant mechanism [*Al-Futaisi and Patzek, 2003; Valvatne and Blunt, 2004*]. The important point is that snap-off and cooperative filling may lead to disconnection and bypassing of the CO₂. The macroscopic consequences of these pore-scale processes are trapping and relative permeability hysteresis. In accordance with the pore-scale explanation give above, experimental data strongly suggest that the nonwetting phase experiences much more pronounced hysteresis than the wetting phase. Several hysteresis empirical models have been developed to characterize the relative permeabilities and trapped saturation of the nonwetting phase after a flow reversal [*Land, 1968; Killough, 1976; Carlson, 1981; Lenhard and Parker, 1987; Jerauld, 1997; Larsen and Skauge, 1998; Lenhard and Oostrom, 1998; Blunt, 2000*].

2.2. Basis for Residual Trapping in Geologic CO₂ Storage

[12] Saline aquifers, predominantly water wet, are prime candidates for geologic CO₂ sequestration. In water wet media and a capillary-dominated flow regime, snap-off is the dominant trapping mechanism at the pore scale. Capillary trapping of the nonwetting gas phase occurs during water flooding when the gas saturation is decreasing, and the water saturation increases as it invades the pore space.

During the injection of CO₂ in the geologic formation, the gas saturation increases in a drainage-like process. Vertical flow paths are created as the gas phase migrates laterally away from the injection wells and to the top of the aquifer due to buoyancy forces. Once the injection stops, the CO₂ continues to migrate upward. At the leading edge of the CO₂ plume, gas continues to displace water in a drainage process (increasing gas saturation), while at the trailing edge water displaces gas in an imbibition process (increasing water saturations). The presence of an imbibition saturation path leads to snap-off and, subsequently, trapping of the gas phase. A trail of residual, immobile CO₂ is left behind the plume as it migrates upward (Figure 1).

3. Numerical Simulations

[13] We perform a series of representative simulations to assess the impact of residual trapping and relative permeability hysteresis on the migration and distribution of injected CO₂ in a sequestration project. We used the commercial reservoir simulator Eclipse 100 [*Schlumberger, 2005*].

3.1. Reservoir Description

[14] We carried out simulations of CO₂ injection in a synthetic but realistic model of a geologic formation. Since the formation is relatively small, these simulations are representative of a CO₂ pilot test, rather than a full-field CO₂ sequestration project. The objective is to illustrate the importance of the physical processes considered, although in section 5 we discuss their implications at a more realistic operational scale. We selected the PUNQ-S3 model, which is a geometrically complex and heterogeneous three-dimensional geologic model originally designed as a test case for oil production forecasting under uncertainty. The original PUNQ-S3 model is described in detail elsewhere [*Floris et al., 2001*], and the model data are publicly available for download (Netherlands Institute of Applied Geosciences, PUNQ case Studies, <http://www.nitg.tno.nl/punq/cases/index.shtml>). We modified the original model slightly to study hysteresis and trapping effects in a CO₂ injection scenario. The modifications were limited to the well locations and flow rates, the fluid properties, the depth of the formation, and the relative permeability tables. The geometry of the model is characterized by a dome in the center and contains five layers of fluvial sand and shale. We set the top of the formation at a depth of 840 m. The average reservoir thickness is about 15 m. The formation is discretized into 19 × 28 × 5 grid blocks, of which 1761 blocks are active. The *x* and *y* dimension of each block is 180 m. The average porosity is 0.2, and the average horizontal permeability is 100 md = 10⁻¹³ m². The permeability anisotropy ratio is about 3. A map of the horizontal permeability is shown in Figure 2.

[15] The aquifer pore volume is approximately 3.6 × 10⁶ m³ with an initial pressure of 90 bar at the top of the structure and a temperature of 40°C. These pressure and temperature conditions ensure that the CO₂ is present as a supercritical fluid. Our model has eight injection wells open to the bottom layer of the aquifer. The injection wells are rate controlled and operate with a constraint in the maximum bottom hole pressure of 160 bar, which is never reached. The well terms are handled in a fully implicit

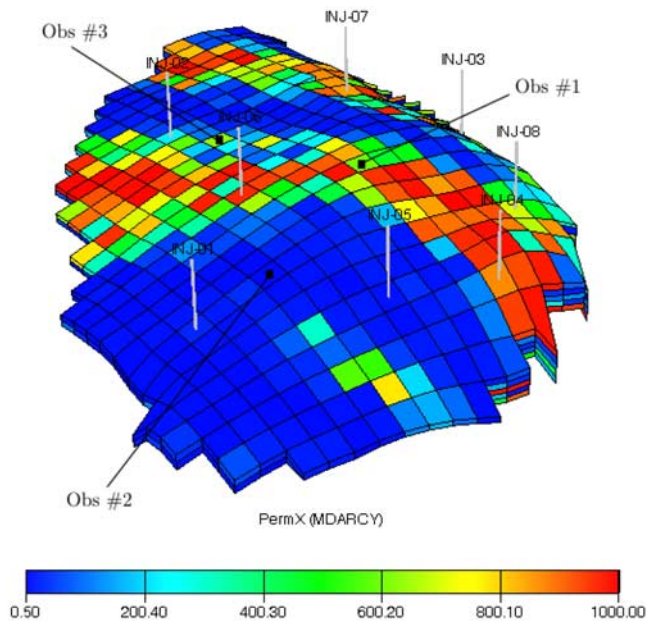


Figure 2. Horizontal permeability distribution (1 md = 10^{-15} m²) of the PUNQ-S3 model. Also shown are the injection wells and the “observation” grid blocks.

fashion, using the default well model in Eclipse. In order to model boundary conditions that are representative of an aquifer that extends beyond the simulation grid, we assign a very large pore volume to the boundary blocks. Specifically, we multiply the pore volume of these blocks by a factor of 1000. While this approach certainly cannot capture the flow dynamics associated with the surrounding aquifer, it has proved to be effective in practice. Upon CO₂ injection, it allows that the brine leave the system, while reproducing a modest pressure buildup at the boundary.

3.2. Trapping and Relative Permeability Hysteresis Models

[16] The relative permeabilities of water and gas are taken from *Oak* [1990] for a water wet Berea sandstone and a gas-water system. Evaluation of the relative permeabilities in Eclipse was performed by linear interpolation between data points. The relative permeability curves are shown in

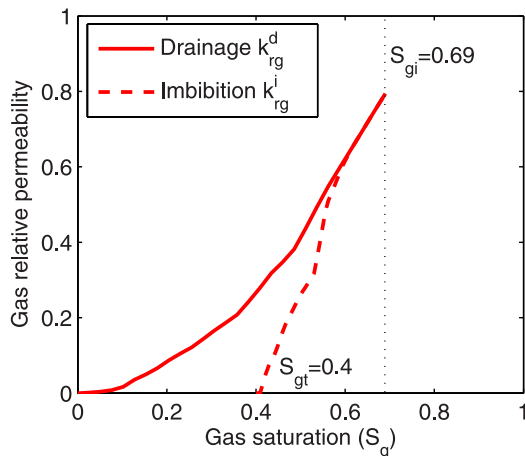


Figure 3. Because the medium is strongly water wet, hysteresis effects are significant in the gas relative permeability only. The relative permeability data of *Oak* [1990] were measured for a brine–nitrogen gas system under low pore pressure and, admittedly, may not be entirely appropriate for a system with supercritical CO₂. The value of residual gas saturation of 40%, however, is in the range of experimental values of *Bennion and Bachu* [2006].

[17] The most important quantity determining the significance of hysteresis effects is the trapped gas saturation after a flow reversal (from drainage to imbibition). A trapping model attempts to relate the trapped (residual) gas saturation to the maximum gas saturation, that is, the actual gas saturation at flow reversal. Most relative permeability hysteresis models make use of the trapping model proposed by *Land* [1968]. In this model, the trapped gas saturation S_{gt} is computed as

$$S_{gt} = \frac{S_{gi}}{1 + CS_{gi}}, \quad (1)$$

where S_{gi} is the initial gas saturation (actual gas saturation at flow reversal) and C is the Land trapping coefficient. The Land trapping coefficient is computed from the bounding drainage and imbibition relative permeability curves as follows:

$$C = \frac{1}{S_{gt,max}} - \frac{1}{S_{g,max}}, \quad (2)$$

where $S_{g,max}$ is the maximum gas saturation, and $S_{gt,max}$ is the maximum trapped saturation, associated with the bounding imbibition curve. All these quantities are illustrated in Figure 4. The Land trapping model has been validated by comparison with experiments [*Land*, 1968; *Jerauld*, 1997; *Spiteri and Juanes*, 2006] and pore-network simulations [*Spiteri et al.*, 2005] for water wet rocks. The bounding drainage and imbibition curves from the experimental data (Figure 3) result in a Land trapping coefficient $C \approx 1$.

[18] A relative permeability hysteresis model characterizes the scanning curves during imbibition and drainage cycles. In this paper, we have used the *Killough* [1976]

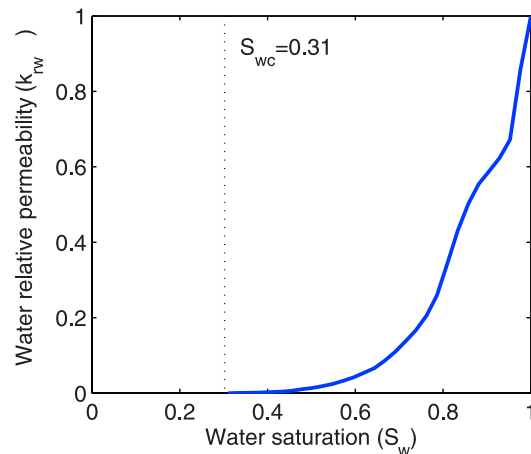


Figure 3. Relative permeability curves used in the CO₂ simulations, taken from *Oak* [1990] for a water wet Berea sandstone.

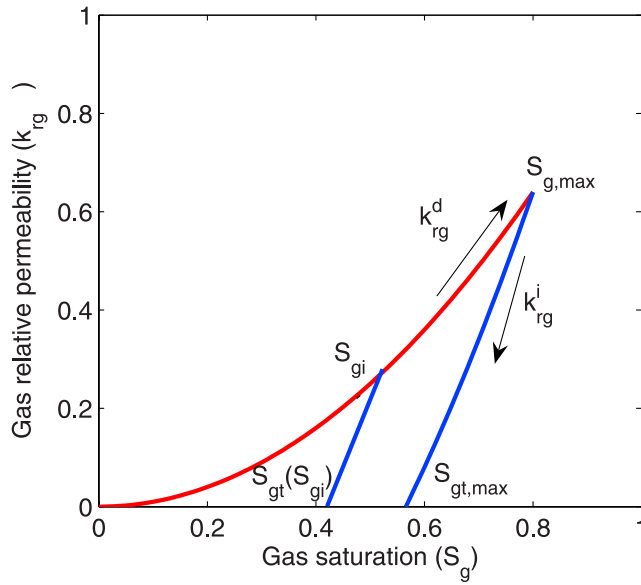


Figure 4. Parameters required in the evaluation and application of the Land trapping model.

hysteresis model. In Killough's method, the gas relative permeability along a scanning curve, such as the one depicted in Figure 4, is computed as:

$$k_{rg}^i(S_g) = k_{rg}^{ib}(S_g^*) \frac{k_{rg}^d(S_{gi})}{k_{rg}^d(S_{gt,max})}, \quad (3)$$

where

$$S_g^* = S_{gt,max} + \frac{(S_g - S_{gt})(S_{gi,max} - S_{gt,max})}{S_{gi} - S_{gt}}. \quad (4)$$

In equation (3), k_{rg}^d and k_{rg}^{ib} represent the bounding drainage and imbibition curves, respectively. The bounding imbibition curve is assumed to be available from experiments, or computed using Land's imbibition model [Land, 1968]. In Killough's model, scanning curves are assumed to be "reversible", so that the imbibition curve is representative of a subsequent drainage process.

[19] Of course, capillary pressure–saturation relationships also exhibit marked hysteresis effects. Several mathematical models exist to treat hysteretic capillary pressure curves, including the one proposed by Killough [1976]. From a practical point of view, however, capillary pressure effects are often negligible at the time of numerically simulating field-scale displacements, especially, as is the case here, when the characteristic capillary length is much smaller than the grid resolution [Aziz and Settari, 1979; Spiteri and Juanes, 2006]. We have indeed checked that our predictions are insensitive to the choice of the hysteretic or nonhysteretic capillary pressure curves. As a result, capillary pressure was assumed to be zero in the simulations presented here.

3.3. Fluid Properties

[20] We have used fluid properties that are representative of water and CO₂ at reservoir conditions [McCain,

1990; Garcia, 2003]. The phase behavior of the system is greatly simplified by the assumption that the two fluids are immiscible.

[21] Assuming a concentration of total dissolved solids of about 5%, we used a value of 1030 kg m⁻³ for the density of water at standard conditions. A constant value of water compressibility 4.35×10^{-5} bar⁻¹ was employed. For a reference pressure of 130 bar and a reference temperature of 40°C, the dynamic viscosity of water is approximately 0.81×10^{-3} kg m⁻¹ s⁻¹. We assumed a linear dependence of water viscosity with pressure, with a viscosibility

$$\frac{1}{\mu_w} \frac{\partial \mu_w}{\partial p} = 7.4 \times 10^{-4} \text{ bar}^{-1}. \quad (5)$$

[22] We relied on the compilation and analysis of Garcia [2003] for the determination of appropriate CO₂ properties at reservoir conditions. From the density at reservoir and surface conditions, one can determine the CO₂ formation volume factor B_g , which is defined as the volume of CO₂ at reservoir conditions corresponding to a unit volume of CO₂ at surface (standard) conditions. Thus

$$B_g = \frac{V_{res}}{V_{sc}} = \frac{\rho_{sc}}{\rho_{res}}. \quad (6)$$

In Table 1 we list the density, formation volume factor and dynamic viscosity of CO₂ at 40°C as a function of pressure. Linear interpolation between values in Table 1 was used in the simulations.

3.4. Setup of Numerical Simulations

[23] The formation is initially filled with brine. CO₂ injection is simulated by controlling the volume of CO₂ that is injected into the formation. A total of 0.15 pore volumes are injected into the bottom layer of the aquifer. The injection rate and volume are the same for all eight injectors.

[24] We simulated four different scenarios, summarized in Table 2. They were designed to assess the following factors: (1) hysteresis and trapping, by comparing the results of case 2 (in which hysteresis is modeled) and case 1 (in which hysteresis is not modeled); (2) injection rate, by comparing the results of case 2 (CO₂ injection during 10 years) and case 3 (injection of the same volume over a period of 50 years); and (3) injection of water, by comparing case 4 (injection of 0.05 pore volumes of water after 5 and 10 years of CO₂ injection) with case 2, respectively (injection of CO₂ alone).

[25] For each case, we show results of the fluid distribution after 500 years from the beginning of the injection phase. We

Table 1. CO₂ Properties at 40°C

Pressure, bars	Density, kg m ⁻³	Formation Volume Factor	Viscosity, $\times 10^{-3}$ kg m ⁻¹ s ⁻¹
1	1.8	1.00000	
80	234	0.00769	0.0325
100	447	0.00403	0.0525
120	632	0.00285	0.0625
160	747	0.00241	0.0725
200	803	0.00224	0.0800
300	883	0.00204	0.0950

Table 2. Summary of CO₂ Sequestration Simulations

Case	Description	Injection Scheme	Injection Time
1	no hysteresis	0.15 PV CO ₂	10 yr
2	hysteresis	0.15 PV CO ₂	10 yr
3	hysteresis	0.15 PV CO ₂	50 yr
4	hysteresis – WAG	0.075 PV CO ₂ , 0.05 PV water 0.075 PV CO ₂ , 0.05 PV water	5 yr, 1 yr 5 yr, 1 yr

plot three-dimensional views of the CO₂ saturation. Additional insight into the dynamic behavior of the system is gained by plotting the evolution of the actual CO₂ saturation and the trapped CO₂ saturation at specific grid blocks. We chose three “observation” points: one near the top of the anticline structure (observation point 1, grid block 13,18,1), one at a slightly lower elevation (observation point 2, grid block 7,21,1), and a third one at a lower elevation still (observation point 3, grid block 11,11,1). All of them are located at the top layer of the formation. The location of the injection wells and the observation grid blocks is shown in Figure 2.

4. Results and Discussions

4.1. Effect of Hysteresis and Trapping

[26] We begin by illustrating the dramatic effect of relative permeability hysteresis on the predictions of the fate of the injected CO₂. We compare the results from case 1 (no hysteresis) and case 2 (with hysteresis). In case 1, the gas relative permeability is assumed to be reversible, and only the drainage curve is used. Both cases simulate injection of a total of 0.15 pore volumes of CO₂ during 10 years, and the migration of the CO₂ plume during the next 490 years.

[27] In Figure 5 we plot the distribution of CO₂ saturation predicted by both models after 500 years from the beginning of injection. In case 1, because the gas relative permeability is assumed to be reversible, the model does not predict any trapping of CO₂. The CO₂ plume migrates upward due to buoyancy forces without leaving any residual saturation behind. After a sufficiently long time, the model predicts the formation of a gas cap of mobile CO₂ at the top of the formation. This scenario is unfavorable from a sequestration standpoint: damage in the cap rock could lead to fractures that might serve as conduits for leaks of the mobile CO₂ to upper formations and, eventually, the atmosphere.

[28] The predictions under case 2 are entirely different. After the injection phase, the model predicts a trail of residual, immobile CO₂ during the migration of the plume. As a result, while there is a net flow of CO₂ in the vertical direction, trapping prevents the injected CO₂ from forming a gas cap. In fact, the simulation predicts that, after 500 years or less, almost all the CO₂ is trapped in the formation. Accounting for hysteresis effects leads to a spread out distribution of trapped CO₂, as opposed to a concentrated distribution of mobile CO₂. This scenario is in fact much more realistic and, importantly, much more favorable for the effectiveness of CO₂ sequestration: it minimizes the risk of leaks (the gas is immobile) and enhances other sequestration mechanisms such as dissolution into the brine and geochemical binding (more interfacial area between the CO₂ and the initial pore water).

[29] Figure 6 shows the evolution of the actual CO₂ saturation with time at the three different observation grid blocks in the aquifer during the first 200 years. For the block at the very top (observation point 1) we see the accumulation of CO₂ when hysteresis effects are ignored, reaching a saturation value close to 0.7 that corresponds to

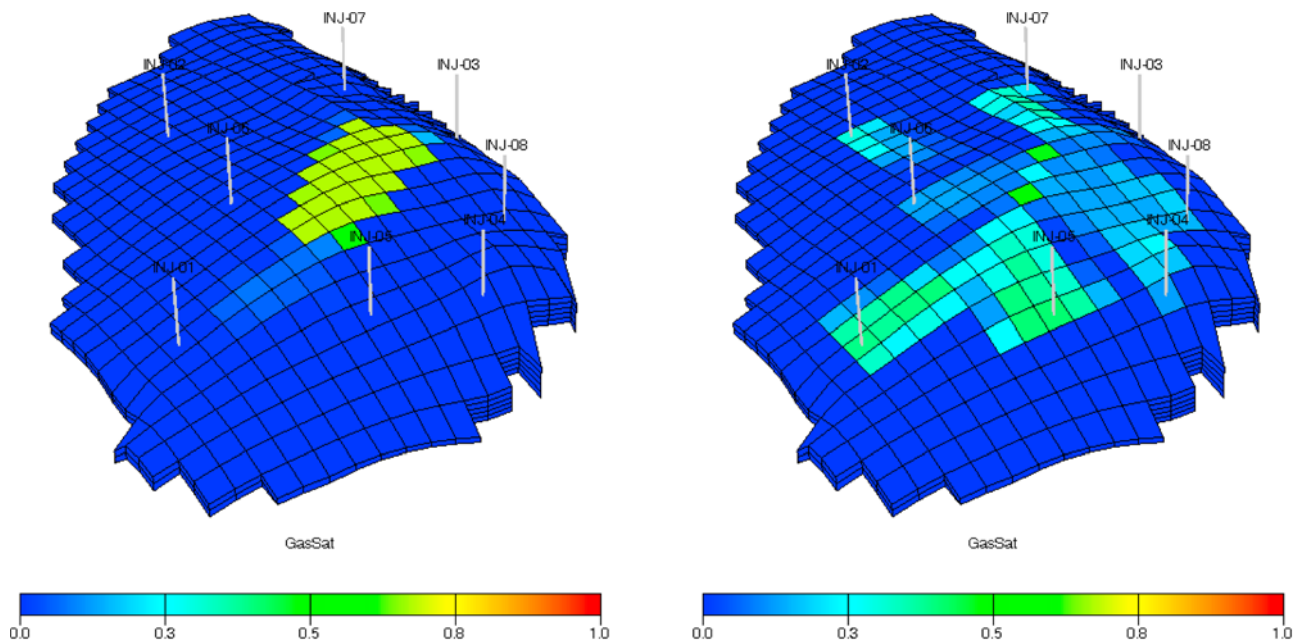


Figure 5. CO₂ saturation distributions after 500 years from the beginning of CO₂ injection. (left) Results from case 1 (no hysteresis). (right) Results from case 2 (with hysteresis).

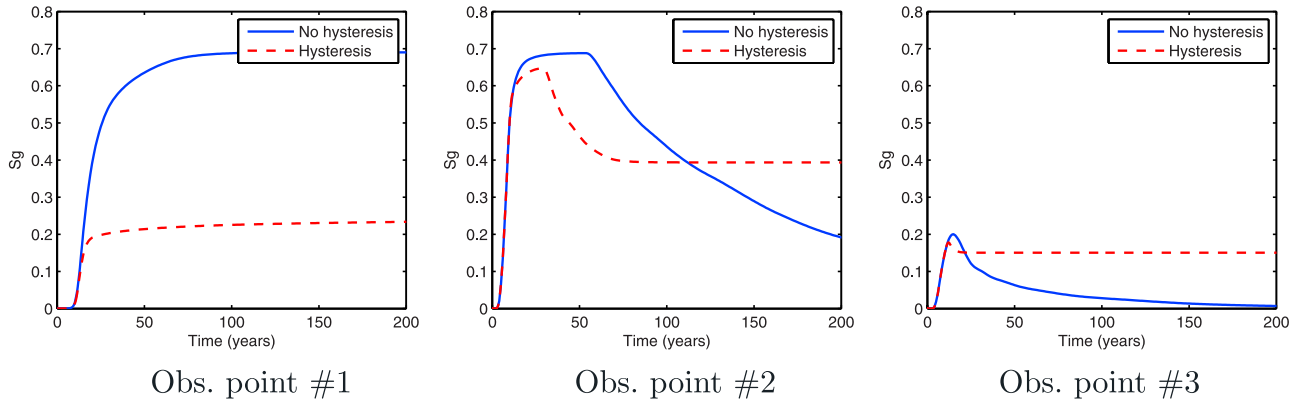


Figure 6. Evolution of actual CO₂ saturations at the three “observation” grid blocks for case 1 (no hysteresis) and case 2 (with hysteresis).

the connate water saturation. In contrast, very little accumulation (CO₂ saturation of about 0.2) occurs when hysteresis is accounted for. The evolution of the CO₂ saturation in blocks at lower elevations (observation points 2 and 3) displays an interesting behavior. When hysteresis is ignored, the gas saturation first increases sharply as the leading edge of the plume reaches the block during its migration upward, and then decreases, eventually to a very low value; the plume travels through the block without leaving any residual CO₂. On the other hand, simulations that account for hysteresis predict that the CO₂ saturation decreases only to a finite, positive value. This is due to trapping during the imbibition process that occurs at the trailing edge of the plume, which results in residual CO₂ being left behind.

[30] These observations are confirmed by the evolution of the trapped CO₂ saturation at the three observation grid blocks, shown in Figure 7. In the model without hysteresis, the trapped CO₂ saturation is of course zero at all times. In the model with hysteresis, the trapped saturation at the top of the anticline (observation point 1) is also equal to zero, because this block only experiences drainage; the actual CO₂ saturation increases with time. At observation points 2 and 3, however, the model predicts that the fraction of trapped CO₂ increases with time. The curve of trapped saturation levels off after 25 years at observation point 3,

which is at a lower elevation, and after almost 100 years at observation point 2 (higher elevation) reflecting the different times at which the trailing edge of the plume travels through the different blocks. In these two grid blocks, all the CO₂ present after 100 years is immobile. The value of the ultimate residual CO₂ saturation is related to the maximum CO₂ saturation through the Land trapping relation, equation (1).

4.2. Effect of Injection Rate

[31] We now investigate the effects of CO₂ injection rate on the overall performance of the sequestration project. We do so by comparing cases 2 and 3, both of which account for hysteresis. We inject the same amount of CO₂ in both cases, but over a period of 10 years in case 2 and over a period of 50 years in case 3.

[32] The saturation distributions for both of these cases are shown in Figure 8. Clearly, a slower injection rate (case 3) leads to more mobile CO₂ reaching the top of the aquifer, which has an adverse effect for sequestration purposes. These results can be explained as follows.

[33] Higher injection rates lead to a more radial displacement pattern and higher gas pressures in the vicinity of wells. Physically, this higher pressure is responsible for the nonwetting gas to invade smaller pores, which have a higher

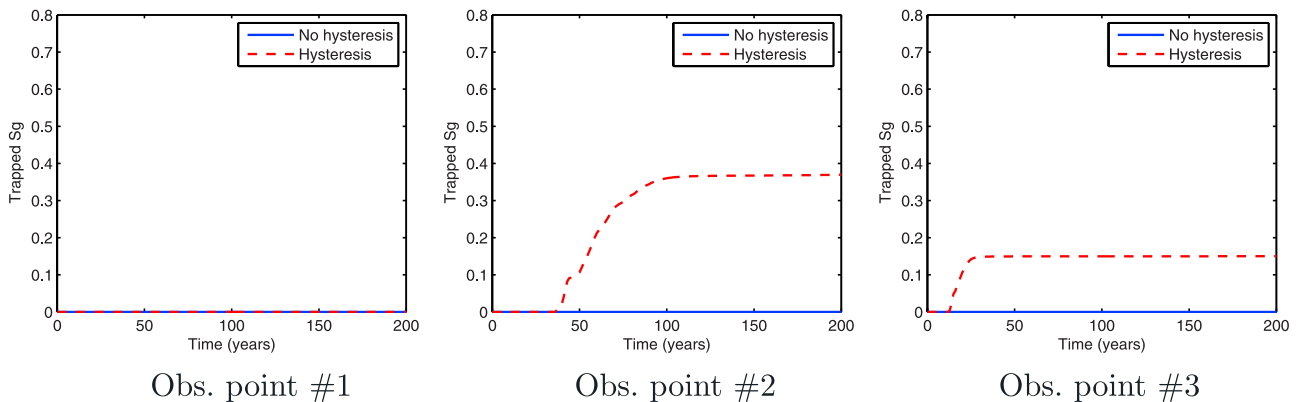


Figure 7. Evolution of trapped CO₂ saturations at the three “observation” grid blocks for case 1 (no hysteresis) and case 2 (with hysteresis).

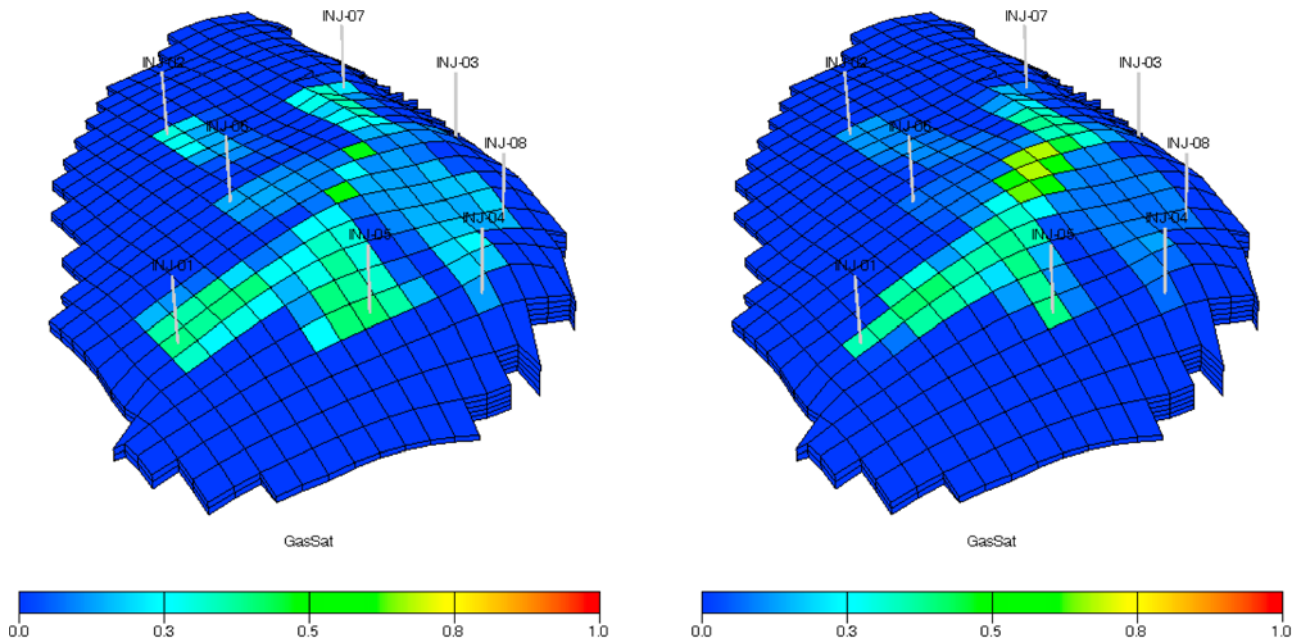


Figure 8. CO₂ saturation distributions after 500 years from the beginning of CO₂ injection. (left) Results from case 2 (injection over 10 years). (right) Results from case 3 (injection over 50 years).

capillary entry pressure. Snap-off occurs readily in smaller pores during imbibition, resulting in increased macroscopic trapping.

[34] For lower injection rates, on the other hand, gravity affects the displacement pattern of water by CO₂ earlier. The low-viscosity gas forms stable paths through the high permeability regions of the porous medium. Microscopically, only the largest pores are invaded, which leads to reduced snap-off during an eventual imbibition process. Moreover, the CO₂ reaches the top of the formation before injection ceases. Since trapping only occurs during imbibition (water displacing CO₂ at the trailing edge of the plume), the fraction of CO₂ that has reached the top of the aquifer cannot undergo imbibition and is therefore not subject to trapping.

[35] The plots of actual and trapped CO₂ saturation over time for these two cases are shown in Figures 9 and 10, respectively. As seen from the plots of observation points 2

and 3, the system undergoes trapping of the CO₂ in both cases. However, the case with low injection rate shows that a higher fraction of CO₂ reaches the top of the anticline (observation point 1): a CO₂ saturation of about 0.7 as opposed to 0.2 of the high-rate injection case.

4.3. Effect of Alternating Water Injection

[36] Finally, we investigate how the performance of the CO₂ sequestration project is affected by the injection of alternating slugs of water and CO₂ into the aquifer. The motivation is to enhance the imbibition process that naturally occurs at the trail of the nonwetting CO₂ plume as it migrates upward.

[37] In Figure 11 we compare the fluid distributions of case 2 and case 4. In both cases we inject 0.15 pore volumes of CO₂ over a period of 10 years. In case 2, the injection of CO₂ is continuous. In case 4, on the other hand, the injection scheme is as follows: 0.075 PV of CO₂ during

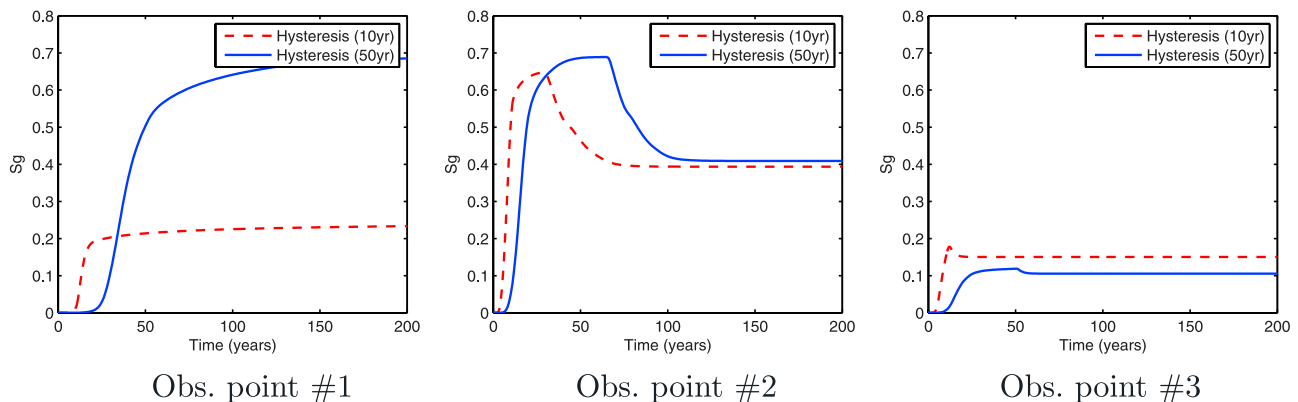


Figure 9. Evolution of actual CO₂ saturations at the three “observation” grid blocks for case 2 (injection over 10 years) and case 3 (injection over 50 years).

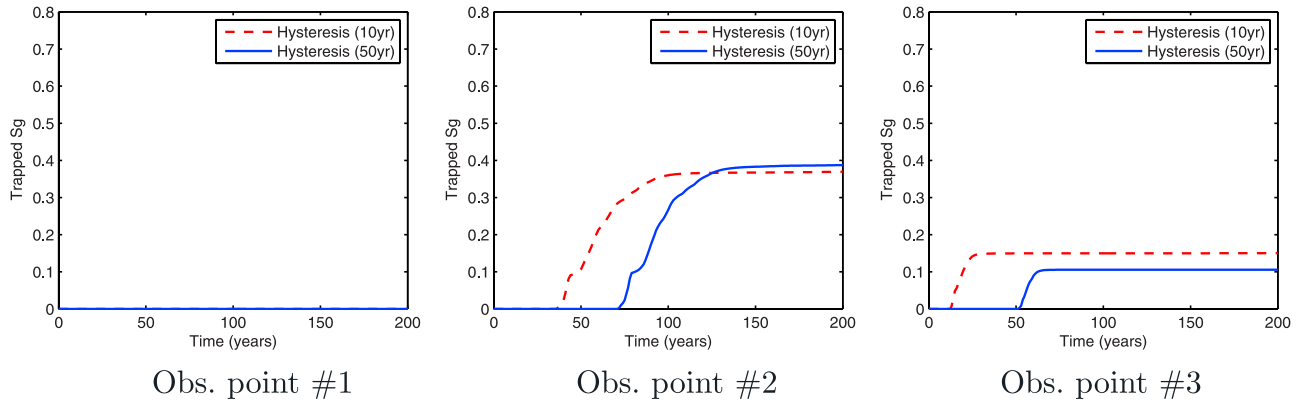


Figure 10. Evolution of trapped CO₂ saturations at the three “observation” grid blocks for case 2 (injection over 10 years) and case 3 (injection over 50 years).

5 years, 0.05 PV of water during one year, 0.075 PV of CO₂ during the next 5 years and, finally, another slug of 0.05 PV of water during one year. The results confirm the expected response: alternating water injection induces more trapping and reduces significantly the amount of CO₂ that accumulates at the top of the aquifer.

[38] The evolution of the actual and trapped CO₂ saturation at the observation grid blocks (Figures 12 and 13, respectively) offers additional insight into the behavior of the displacement process. During the CO₂ injection period, the first 5 years, the curves corresponding to case 2 and case 4 are of course identical. After 5 years, the curves start to deviate due to the injection of water. Water displaces CO₂ radially away from the wells, which can be seen as a forced imbibition process that leads to enhanced trapping. It

explains why less CO₂ reaches the top of the aquifer (observation point 1): gas saturation of about 0.1 instead of 0.2 of the continuous CO₂ injection case.

4.4. Impact on the Bottom Hole Pressure

[39] In the previous sections, we illustrated the benefit of higher injection rates and water-alternating-gas injection in terms of increasing residual trapping of CO₂. However, both strategies lead to an increase in the bottom hole pressure (BHP) at the wells, if injection rate is to be maintained.

[40] In this section, we compare the BHP at two of the wells for case 2 (continuous CO₂ injection) and case 4 (water-alternating-gas injection). The evolution of the BHP at two wells during the injection period is shown in Figure 14. The two wells shown, wells 2 and 4, were selected because

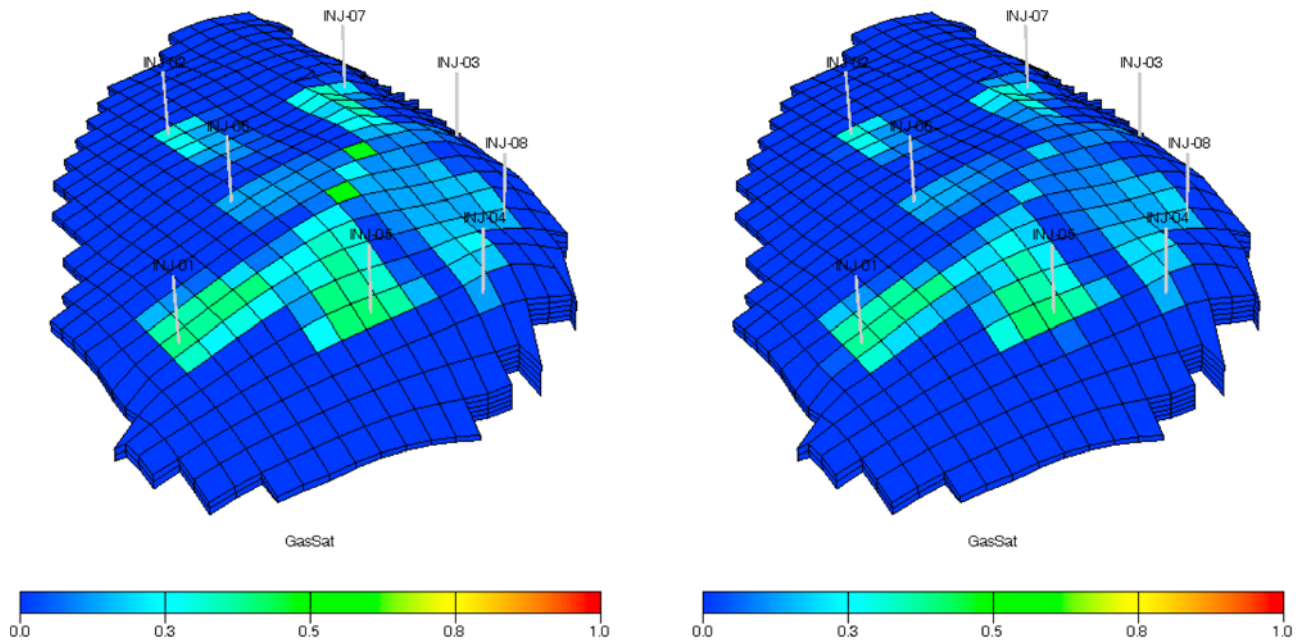


Figure 11. CO₂ saturation distributions after 500 years from the beginning of CO₂ injection. (left) Results from case 2 (no water injection). (right) Results from case 4 (alternating water and CO₂ injection).

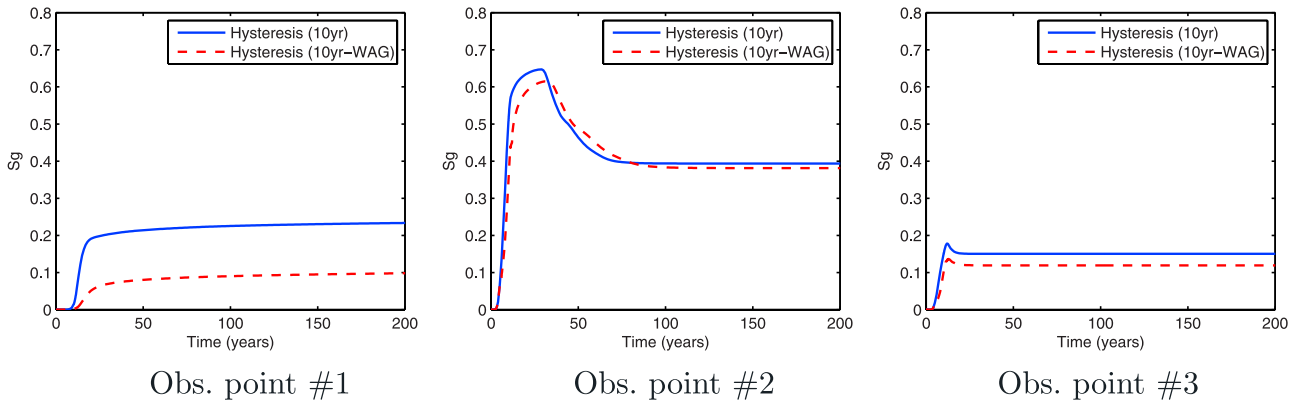


Figure 12. Evolution of actual CO₂ saturations at the three “observation” grid blocks for case 2 (no water injection) and case 4 (alternating water and CO₂ injection).

they display the smallest and largest differences in BHP between the two injection schemes, respectively. It is apparent that a water-alternating-gas strategy results in a significant increase in the operating BHP. The reason is twofold: first, the compressibility of water is much smaller than that of CO₂; moreover, the water injection rate is about three times higher than the injection rate of CO₂. The combined effect is a sudden jump in BHP when water injection starts, and a steeper rate of increase of BHP during water injection. As soon as water injection ceases, the well BHP drops quickly.

[41] The importance of such difference in bottom hole pressure is that BHP is often subject to regulatory limits [Bachu *et al.*, 2003], and a higher BHP also leads to higher operating costs. These factors must be taken into account when establishing the operational benefit of a water-alternating-gas injection strategy.

4.5. Impact of Grid Refinement

[42] The simulations presented so far employ a very coarse discretization (each grid block is approximately 180 × 180 × 3 m³). Clearly, this crude spatial discretization affects the accuracy of the predictions. In particular, the averaging (or mixing) that takes place at each grid block overestimates the exposure of CO₂ to the pore space during

the injection period, thereby overestimating the amount that is trapped during the subsequent imbibition process [Mo *et al.*, 2005].

[43] To illustrate this point, we repeated the simulations using a refined grid in which each grid block is divided into 3 × 3 × 3 grid blocks. This refinement results in a model of about 50,000 active grid blocks. The CO₂ saturation maps after 500 years for case 1 (no hysteresis) and case 2 (with hysteresis) are shown in Figure 15. We show the three-dimensional model up to a vertical cross section, to visualize the vertical CO₂ distribution. When compared with the coarse model in Figure 5 we observe that, for case 2, a higher fraction of CO₂ has reached the top of the anticline. The reason is that a finer discretization captures the overall sweep more accurately, which is overestimated in coarse models. This is particularly relevant here, since the buoyant CO₂ migrates in a relatively thin layer below the reservoir top.

[44] This behavior is also reflected in the evolution of the CO₂ saturation at the observations points (Figure 16). In the refined model, observation point 3 is never contacted by CO₂. At observation point 1, both models (with and without hysteresis) predict the same saturation history, because the CO₂ makes its way to the top of the anticline during (or

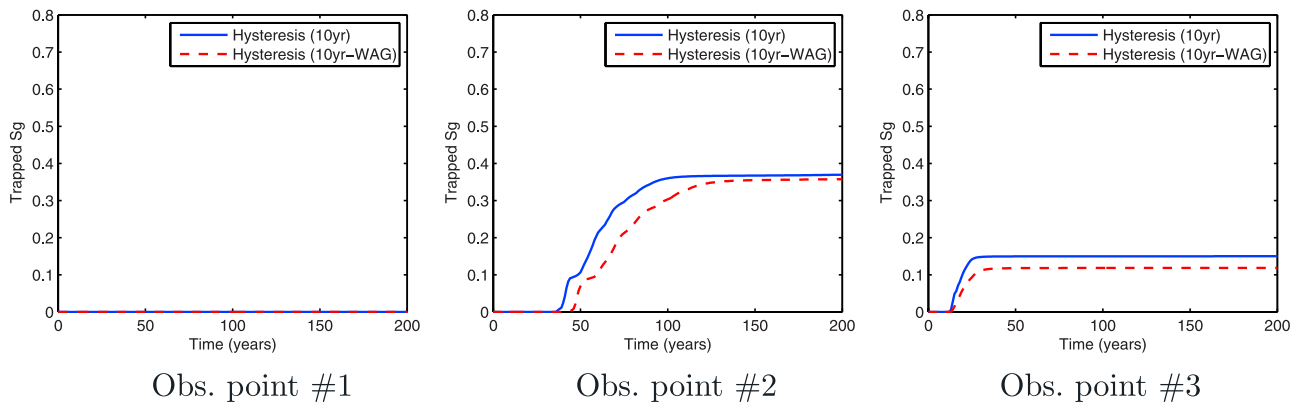


Figure 13. Evolution of trapped CO₂ saturations at the three “observation” grid blocks for case 2 (no water injection) and case 4 (alternating water and CO₂ injection).

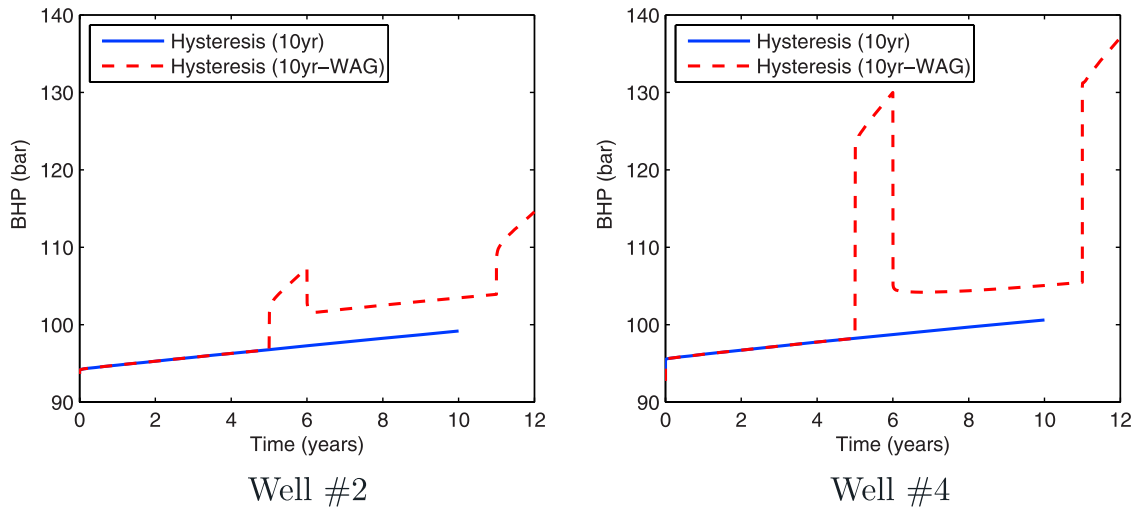


Figure 14. Evolution of bottom hole pressure (BHP) at two different wells for case 2 (continuous CO₂ injection) and case 4 (water-alternating-gas injection).

shortly after) the injection period, from which it cannot be trapped. For observation point 2, however, the simulation shows that wherever the CO₂ plume contacts the reservoir, it will leave a trail of residual immobile gas during its lateral and upward migration.

5. Summary and Conclusions

[45] In this paper, we have investigated the impact of trapping and relative permeability hysteresis in the context of CO₂ sequestration projects in saline aquifers. The main findings can be summarized as follows.

[46] 1. Accounting for trapping and relative permeability hysteresis of the nonwetting CO₂ phase is essential in order to correctly characterize the migration and final distribution of the injected CO₂. Trapping of the CO₂ occurs during the

upward migration of the CO₂ plume, but only after injection has stopped and the trailing edge of the plume is naturally being displaced by water. This imbibition process leads to trapping of the CO₂. A trail of residual CO₂ is left behind as the plume migrates upward.

[47] 2. Trapping of the CO₂ leads to more favorable scenarios for sequestration purposes: a large fraction of the CO₂ is trapped and immobile for practical purposes, and is more spread out throughout the aquifer, thereby increasing the interfacial area for subsequent dissolution in the brine. The simulations presented here are representative of a pilot CO₂ injection test, rather than a full-field storage project. Larger volumes per injection well would result in a CO₂ plume that is less distributed across the reservoir, which would result in less capillary trapping. In contrast, target storage formations would likely be thicker,

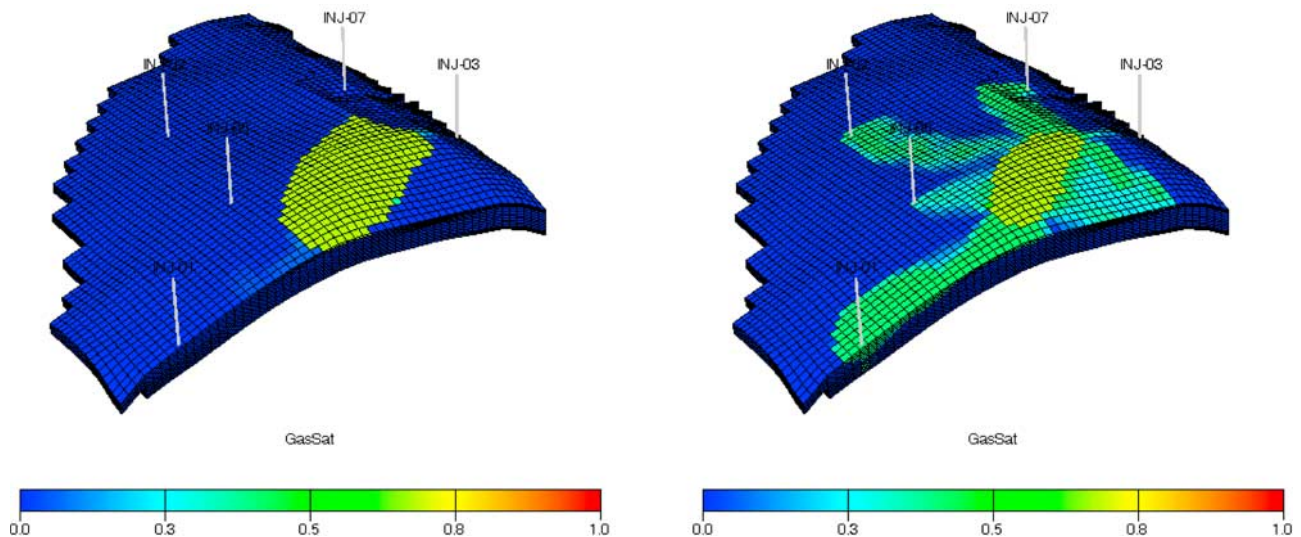


Figure 15. CO₂ saturation distributions after 500 years from the beginning of CO₂ injection, computed on the refined grid. (left) Results from case 1 (no hysteresis). (right) Results from case 2 (with hysteresis).

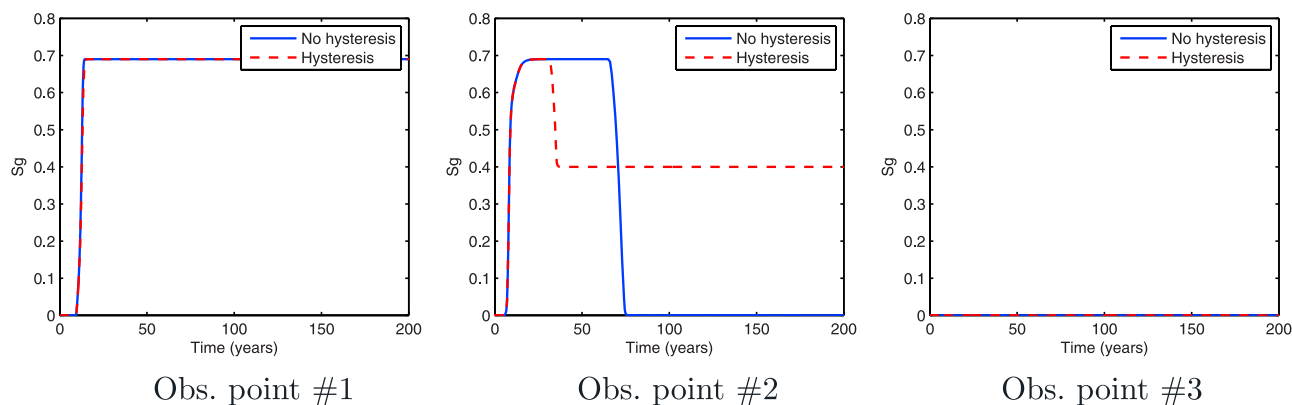


Figure 16. Evolution of actual CO₂ saturations at the three “observation” grid blocks for case 1 (no hysteresis) and case 2 (with hysteresis) from the fine-grid simulation.

sloping aquifers, where the injected CO₂ would continue to rise leaving a trail of residual CO₂ over a larger areal extension.

[48] 3. High injection rates result in more effective sequestration of the CO₂. A shorter injection period leaves less time for the buoyant CO₂ to reach the top of the formation, from which it is difficult to immobilize.

[49] 4. Injection of water slugs alternating CO₂ injection (in the spirit of classical WAG for enhanced oil recovery [Spiteri and Juanes, 2006]) increases the effectiveness of the sequestration project. The injected water forces breakup of large connected CO₂ plumes, enhancing trapping and immobilization of the CO₂. On the other hand, a WAG strategy leads to higher bottom hole pressures at injection wells, which may be limited by seal integrity, regulatory or economical constraints. The identification of WAG as a potentially effective strategy for CO₂ storage lends itself to an optimization problem to maximize the amount of trapped CO₂ by varying the well rates and well completions, subject to BHP constraints.

[50] 5. Coarse simulation models overestimate the sweep and subsequent capillary trapping of CO₂. We believe that an accurate assessment of the different storage mechanisms (hydrodynamic, capillary, solution and mineral trapping) requires high-resolution models that capture the migration paths of the injected CO₂ in the subsurface.

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