# Patchy particles at a hard wall: Orientation-dependent bonding

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- The bulk phase behaviour of this model has been extensively studied by both simulation and theory mostly Werteim's thermodynamic perturbation theory (TPT).
- But TPT cannot account for the orienting effects of surfaces [e.g., Gnan *et al.*, J. Chem. Phys. **137**, 084704 (2012)] it averages over the directionality of interpatch interactions. How do we fix this?

#### The model we simulate

• The interparticle pair potential is the sum of HS repulsion between cores and attraction between surface patches:

$$u(\mathbf{r}_i, \hat{\mathbf{r}}_j) = u_{HS}(\mathbf{r}_{ij}) + \sum_{\alpha, \beta=1}^2 V_{ij, \alpha\beta}$$

• Bol-Kern-Frenkel potential:

$$V_{12,\alpha\beta}(\mathbf{r}_{12}, \hat{\mathbf{r}}_{\alpha 1}, \hat{\mathbf{r}}_{\beta 2}) = V_{\alpha\beta}^{SW}(r_{12})G(\hat{\mathbf{r}}_{12}, \hat{\mathbf{r}}_{\alpha 1}, \hat{\mathbf{r}}_{\beta 2})$$

$$G(\hat{\mathbf{r}}_{12}, \hat{\mathbf{r}}_{\alpha 1}, \hat{\mathbf{r}}_{\beta 2}) = \begin{cases} 1 & \left\{ \text{if } \begin{array}{c} \hat{\mathbf{r}}_{12} \cdot \hat{\mathbf{r}}_{\alpha 1} > \cos\theta_{\alpha\beta}^{max} \\ 0 & -\hat{\mathbf{r}}_{12} \cdot \hat{\mathbf{r}}_{\beta 2} > \cos\theta_{\alpha\beta}^{max} \\ 0 & \text{otherwise.} \end{array} \right.$$

$$V_{\alpha\beta}^{SW}(x) = \begin{cases} \infty & \text{if } x < \sigma \\ -\epsilon_{\alpha\beta} & \text{if } \sigma < x < \sigma + \delta_{\alpha\beta} \\ 0 & \text{otherwise,} \end{cases}$$

δ/2 ....

 $r_{12}$ 

# How we fix the theory

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- Particles with one patch at either pole want to lie parallel to an impenetrable wall.
- Do the theory as if the model contained an effective orientationdependent potential that favours pole-to.pole alignment.

$$V_{eff}(\mathbf{r}_{ij},\omega_i,\omega_j) = v(202) + v(022) + v(222)$$



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#### The fix in more detail

• These are spherical harmonic expansion coefficients:

$$v(l_1 l_2 l) = \sum_{m_1, m_2, m} v(l_1 l_2 l; r_{ij}) C(l_1 l_2 l : m_1 m_2 m) Y_{l_1 m_1}(\omega_i) Y_{l_2 m_2}(\omega_j) Y_{lm}^*(\omega_{ij})$$

• Use truncated and shifted generalised Lennard-Jones r-dependence:

$$v(l_{1}l_{2}l; r_{ij}) = \begin{cases} \infty & \text{if } r_{ij} < \sigma \\ \epsilon_{l_{1}l_{2}l} \left[ \left( \frac{\sigma}{r_{ij}} \right)^{24} - \left( \frac{\sigma}{r_{ij}} \right)^{n} \right] - \epsilon_{l_{1}l_{2}l} \left[ \left( \frac{\sigma}{r_{max}} \right)^{24} - \left( \frac{\sigma}{r_{max}} \right)^{n} \right] & \text{if } \sigma \le r_{ij} < r_{max} \end{cases}$$

Potential parameters used:

$\cos \theta^{max}$	$\delta/\sigma$	$\epsilon_{202}/\varepsilon = \epsilon_{022}/\varepsilon$	$\epsilon_{222}/\varepsilon$	n
0.895	0.119	0.6	0.6	4

Only free parameter is cutoff ("range")  $r_{max}$ .

Helmholtz free energy (FE) functional is sum of three contributions:

$$\mathcal{F}\left[\rho(z), \hat{f}(z, \theta)\right] = \mathcal{F}_{id}\left[\rho(z), \hat{f}(z, \theta)\right] + \mathcal{F}_{hs+b}\left[\rho(z), X(z)\right] + \mathcal{F}_{MF}\left[\rho(z), \hat{f}(z, \theta)\right]$$

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• Translational + rotational entropy of ideal system (exact):

$$\mathcal{F}_{id}\left[\rho(\mathbf{r})\right] = k_B T \int d\mathbf{r} \,\rho(\mathbf{r}) \left\{ \ln\left[\Lambda^3 \rho(\mathbf{r})\right] - 1 \right\} + k_B T \int dz \,d\omega \,\rho(z) \hat{f}(z,\theta) \log\left[4\pi \hat{f}(z,\theta)\right]$$

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• Effective potential contribution, in MF approximation:

$$\mathcal{F}_{MF}\left[\rho(z),\eta(z)\right] = \frac{1}{2} \int dz_i dz_j \,\rho(z_i) \left[\overline{v}(202;|z_i - z_j|)\eta(z_i) + \overline{v}(022;|z_i - z_j|)\eta(z_j)\right] \rho(z_j) \\ + \frac{1}{2} \int dz_i dz_j \,\rho(z_i)\eta(z_i)\overline{v}(222;|z_i - z_j|)\rho(z_j)\eta(z_j)$$

• Minimise grand-canonical potential:

$$\begin{aligned} \frac{\Omega\left[\rho(z),\hat{f}(z,\theta)\right]}{A} &= \mathcal{F}\left[\rho(z),\hat{f}(z,\theta)\right] + \int dz \left[V_{\text{ext}}(z) - \mu\right]\rho(z) \\ \frac{\delta\Omega\left[\rho(z),\hat{f}(z,\theta)\right]}{\delta\rho(z)} &= 0 \Leftrightarrow \frac{\delta\mathcal{F}\left[\rho(z),\hat{f}(z,\theta)\right]}{\delta\rho(z)} = \mu - V_{\text{ext}}(z) \\ \frac{\delta\Omega\left[\rho(z),\hat{f}(z,\theta)\right]}{\delta\hat{f}(z,\theta)} &= \lambda \end{aligned}$$

• Outputs are number density profile  $\rho(z)$ , fraction of unbonded sites X(z), and orientational order parameter profile  $\eta(z)$ :

$$\eta(z) = \int P_2(\cos\theta)\hat{f}(z,\theta)\,d\omega$$

- Standard *NVT* MC of N = 10800 particles enclosed in a cubic box of edge length  $30\sigma$  ( $\rho_{bulk}\sigma^3 = 0.4$ ).
- Periodic boundary conditions along directions x and y, two hard walls at  $z = \pm 15\sigma$ .
- Elementary roto-translational moves consist of a random translation of at most  $\pm 0.05\sigma$  and a random rotation of at most  $\pm 0.1$  rad.
- Temperatures from  $k_B T/\varepsilon = 0.30$  down to  $k_B T/\varepsilon = 0.08$ .
- Number of MC steps (one step = N attempts to move a particle) increased progressively from  $10^6$  to  $10^7$  on lowering the temperature. No difficulty equilibrating.

## Results: density and order parameter profiles, $T^* = 0.3, 0.2$



P. I. C. Teixeira and F. Sciortino, J. Chem. Phys. 151, 174903 (2019)

#### Results: density and order parameter profiles, $T^* = 0.15$



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### Results: density and order parameter profiles, $T^* = 0.12$



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### Results: density and order parameter profiles, $T^* = 0.11$



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#### Results: profiles of fraction of unbonded patches



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#### Results: r<sub>max</sub> and mean chain length vs T



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- The theory is unable to predict the weaker bonding at the wall,
- Other patch configurations could be modelled.

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# Shameless publicity plug



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